The NASA Space Radiation Laboratory at Brookhaven National Laboratory

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Outline

The NASA/BNL Space Radiation Program Sampling of experiments The NSRL Facility Operations Beam Characteristics Uniform Beams Dosimetry Large Dynamic Range Camera Imaging System **Solar Particle Simulator Summary**



"Because astronauts are spending more time in space, the National Aeronautics and Space Administration (NASA) is working with Brookhaven National Laboratory and others here on Earth to learn about the possible risks to human beings exposed to space radiation. To study the radiobiological effects using beams that simulate the cosmic rays found in space, a new, \$34-million NASA Space Radiation Laboratory (NSRL) has been established at Brookhaven Lab."

http://www.bnl.gov/medical/NASA/NSRL_description.asp



Human Environments in Space

Present Future 30 day mission, 2015-2020 30 month mission, 1 year mission, 2025-2030 2005



Radiation Doses on Earth and in Space

1y in Houston 1y in Denver 1y in Kerala, India Apollo 14 Skylab 4 Shuttle mission 41-C Mission to Mars at solar minimum

100 mrem 200 mrem 1,300 mrem 1,100 mrem (9-day to the Moon) **18,000 mrem (87-day in orbit)** 5,600 mrem (18-day in orbit) **130,000 mrem (30 month) 30,000 mrem in 1.5 y on Mars** 80,000 mrem in 1 y in space, +20,000 mrem from a solar flare

Chest X-ray PET scan Treatment of brain cancer 50 mrem 1,000 mrem 500,000,000 mrem (to normal brain)



NSRL became operational during summer 2003

> 100 experimenters from 24 institutions (U.S. and abroad)

Brookhaven researchers and other NASA-sponsored scientists irradiate a variety of biological specimens, tissues, and cells, as well as DNA in solution. Other experimenters use industrial materials as samples, studying their suitability for space suits and spacecraft shielding.



Sampling of experiments

- Effect of Deep Space radiation on Human Hematopoietic Stem Cells
- Risk Assessment and chemoprevintion of HZE-induced CNS damage
- Heavy Ion Induced Chromosome Damage and Biomedical Countermeasures
- DNA damage clusters in low level radiation responses of human cells.
- Complex Space Radiation-induced DNA damage Clusters in Human Cell Transformation: Mechanisms, relationships and Mitigation.
- Induction of Bystander Effects by High LET Radiation in Cells
- Gene Expression in the Nematode C. elegans following Irradiation with Charged Particles
- Heavy Ion Particle Impact on Simulated Martian Regolith
- MSL/RAD Technology Demonstration Model Characterization
- Spacecraft shielding and components experiments
- Ion fragmentation experiments



1 GeV Fe tracks in cells.





Local Polyethylene Shielding Study



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Heavy Ions

Charged Ionizing Radiation Outer-Space is full of them



Fig. 3.5. Typical energy spectra for protons, helium ions, carbon ions, and iron ions from "top to bottom," respectively, at solar minimum. The solid line is the local interstellar spectrum (Simpson, 1983a).







Booster Parameters

POWER DIPOLE BACKLED WINDINGS	Parameter	Value
Constant E SUPERPENCO	Circumference	201.78 (1/4 AGS) m
DG STRAIGHT SECTION REMOVE BEAM DUMP & WALL CURPENT MONITOR INSTALL SUPPERT MONITOR	Ave. Radius	32.114 m
NODEY DO 1/4 CELL CHANGER	Magnetic Bend R	13.8656 m
	Lattice Type	Separated Function, FODO
	No. Superperiods	6
	No. of Cells	24
	Betatron Tunes,X,Y	4.82, 4.83
	Vacuum Chamber	70 x 152 mm Dipoles 152 mm (circular) Quads
	Max. Rigidity	17 Tm
BLIP BLIP	Injection Rigidity	2.2 Tm (200 MeV protons)
BROOKHAVEN	Acceleration Rate	8.9 T/s (7.5 Hz)
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Instrumentation Layout







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Operations Beams delivered, NSRL-5 to NSRL-7

Ion	Energy (MeV/n)	Intensity
р	1000	3.4×10 ¹⁰
С	290	$1.2 imes 10^{10}$
0	1000, 600	4.0×10 ⁹
Si	600, 300	3.0×10 ⁹
Cl	500	2.0×10 ⁹
Ti	1100	8.0×10 ⁸
Fe	1000, 600, 300	2.0×10 ⁹





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Extraction Point







300 MeV/n Carbon

2 GeV Protons



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Beam Spills Delivered





Uniform Beams









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256-Element Ion-Chamber (Beam Imaging)

Beam Image on the 256-element IC During E.Blakely's experiment (NSRL-2).





32-Element Ion-Chamber (Dosimetry)















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Solar Particle Simulator

The main motivation is to be able to reproduce energy spectra of the environment in space, particularly solar events. NSRL Ground rules

- Need to be able to irradiate a single sample over entire energy spectrum within as little as a 1 hour period.
- Need a clean beam, as clean as any current NSRL experiment now receives.
- Need to know the energy of the beam as well as any current NSRL experiment.
- Need to know the dose as well as any current NSRL experiment.



Solar Proton Events



Large solar proton event integral fluence spectra at 1 AU.



Actively Changing Energy at NSRL



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Current Capabilities

- 1. Beam energies from 0.05 to 1 GeV/n with any ion that can be produced by a Tandem Van de Graff
- 2. Lowest intensities operated ~10² ions/cm²/cycle
- 3. Smallest beams around 1 cm, smaller possible (not uniform, Gaussian). Largest beams around 20x20 cm² uniform
- 4. Mixed field of protons and ions on a single target
- 5. 1 3 msec pulsed beams
- Fast Extracted Beams (1 to 3 200 nsec pulses in 5 usec, or one 4 usec pulse)
- 7. Solar particle simulator
 - 1. Large range of ion energies over single irradiation
 - 2. Fast energy change



Future Capabilities

EBIS (electron beam ion source) will provide

- All ions up to U, including noble gases
- Higher intensities for current ion set
- Multiple mixed field ions (more than 2 ions species/ irradiation)
- Pulsed synchronized beams (with experimenter signals)
 - ^{1.} Synchronize beam with breathing, heart rate, EKG,
 - 2. For low energy beams, allow for time of flight analysis



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Supplemental



Transverse fields in a normal sextupole are :

$$B_x(x, y) = -6B_s xy$$

$$B_y(x, y) = -3B_s(x^2 - y^2)$$

where,

$$B_s = -\frac{1}{6} \left(\frac{d^2 B_y}{dx^2} \right)_0$$

A particle with a magnetic rigidity $B\rho$ receives (thin lens) kicks by a sextupole of length *L*,

$$\Delta x' = \frac{1}{2} \frac{L}{B\rho} \left(\frac{d^2 B_y}{dx^2} \right)_0 (x^2 - y^2)$$

and,

$$\Delta y' = -\frac{L}{B\rho} \left(\frac{d^2 B_y}{dx^2} \right)_0 (xy)$$



define the normalized sextupole strength as

$$S = \frac{1}{2} \beta^{\frac{3}{2}} \frac{L}{B\rho} \left(\frac{d^2 B_y}{dx^2} \right)_0$$

The Kobayashi Hamiltonian

$$H = \frac{\varepsilon}{2} \left(X^2 + X'^2 \right) + \frac{S}{4} \left(3XX'^2 - X^3 \right)$$

where

$$X \equiv \frac{x}{\sqrt{\beta_x}}$$
 and $X' \equiv \frac{\alpha_x}{\sqrt{\beta_x}} x + \sqrt{\beta_x} x'$

are the normalized phase space coordinates, and,

$$\varepsilon = 6\pi \delta Q = 6\pi (Q_{particle} - Q_{resonance})$$

The first term in H describes particle motion in the linear unperturbed lattice (S = 0). The trajectories are circles with radius $\sqrt{2H/\varepsilon}$ in normalized phase space.

The second term (perturbative term) distorts the circle.

Slow Extraction Dynamics

when H has the value $\left[\left(2\varepsilon/3\right)^3/S^2\right]$ it factors into 3 terms,

$$\left(\frac{S}{4}X + \frac{\varepsilon}{6}\right)\left(\sqrt{3}X' + X - \frac{4\varepsilon}{3S}\right)\left(\sqrt{3}X' - X + \frac{4\varepsilon}{3S}\right) = 0$$

The 3 lines define the boundries between stable and unstable regions of phase space. The size of the stable region is determined by the ratio ε / S .





$$h = \frac{2}{3}\frac{\varepsilon}{S} = \frac{4\pi}{S}\partial Q$$

The area of the stable region is

$$\mathbf{A} = 3\sqrt{3}\boldsymbol{h}^2 = \frac{48\sqrt{3}\pi}{\boldsymbol{S}^2} \left(\partial \boldsymbol{Q}^2\right)\pi$$

The area within a particular particles linear unperturbed motion is called the single particle emittance.

$$E = a^2 \pi$$
, where $a^2 = X_0^2 + X_0^2$

Particle motion remains stable as long as the particle motion lies within the stable triangle.

$$\boldsymbol{E}_{stable} \leq \frac{48\sqrt{3\pi}}{\boldsymbol{S}^2} \left(\delta \boldsymbol{Q}^2 \right) \pi$$

This criteria for stability can be rewritten in terms of absolute betatron tune, with

$$\partial Q = Q_{particle} - Q_{resonance}$$
$$Q_{resonance} - \sqrt{\frac{1}{48\sqrt{3}\pi}} S_{\sqrt{\frac{E}{\pi}}} < Q_{particle} < Q_{resonance} + \sqrt{\frac{1}{48\sqrt{3}\pi}} S_{\sqrt{\frac{E}{\pi}}}$$

This can be shown graphically by plotting the action variable, \sqrt{E} , as a function of betatron tune.

Slow Extraction Dynamics



Extraction Methods:

- 1. Widen stopband by increasing S
- 2. Move particles into resonance by changing betatron tune (AGS).
- 3. Increase particle amplitudes until it encounters the unstable region (rf knockout method).



Booster Resonant/Slow Extraction



Injection, Acceleration, and Tune Space Manipulations







Energy Measurement in an Accelerator

<u>Energy</u>	Symbols:					
β = (2 π R _{booster} f _{rf})/ (h c)	h = harmonic #					
$\gamma = 1/(1-\beta^2)^{1/2}$	c = 2.99792458x10 ⁸ m/s					
K.E. = $(\gamma - 1)$ mc²	R _{booster} = R = 32.113 m					
<u>Uncertainty</u>	f _{rf} = accelerating frequencyValues used in figure:					
$d\beta^2 = (2\pi f_{rf} / hc)^2 dR^2 + (2\pi R/hc)^2 df_{rf}^2$						
$d(K.E.) = mc^2 \beta (1-\beta^2)^{-3/2} d\beta$	dR = 5 mm df _{rf} = 10 Hz h = 3					



Time Structure Formalism

definitions:

Q = horizontal betatron tune ξ = horizontal chromaticity = $\frac{dQ/Q}{dp/p}$ I_m = current in the Main Dipoles and Quadrupoles N = number of particles ($\frac{dN}{dQ}$ represents the particle distribution in tune space

T = period over which particles are extracted

Low frequency duty factor:

$$D_f = rac{[f_T S(t)dt]^2}{T f_T [S(t)]^2 dt} \ \Rightarrow (rac{f_{av}}{f_{rms}})^2 \ in \ general$$

where

$$S(t) = \frac{dN}{dt} = \frac{dN}{dQ}\frac{dQ}{dt}$$

if there is no ripple,

$$S(t) = \frac{dN}{dQ}\dot{Q_0}$$

where \dot{Q}_0 is the rate at which particles move into resonance.

$$\dot{Q_0} = \frac{Q\xi}{I_m} \frac{dI_m}{dt}$$



If there is ripple on the magnet power supplies;

$$S(t)=\;\frac{dN}{dQ}(\dot{Q_0}+\dot{Q_v})\;$$

where Q_v is the variations in the rate at which particles move into resonance.

$$\dot{Q_v} = \frac{Q\xi}{I_m L_m} \sum_h V_h$$

 L_m is the total inductance of the main dipoles and quads and V_h is the sum of the 60 Hz harmonics amplitudes (in volts). Reducing Time structure using RF Phase Displacement

$$S(t) = \frac{dN}{dQ} \dot{Q}_0 (1 + \frac{\dot{Q}_v}{\dot{Q}_0})$$

For 1 particular frequency we can write the duty factor as

$$D_f = \frac{1}{1 + \frac{1}{2}(\frac{\dot{Q}_v}{\dot{Q}_0})^2} = \frac{1}{1 + \frac{1}{2}(\frac{\omega\delta Q}{v_0})^2}$$

where

 $\omega = \text{frequency}$ $\delta Q = \text{relative ampl. of that freq. in tune space}$ $v_0 = \text{speed that beam crosses resonance}$

$$v_0 = \frac{\Delta p}{p} \frac{1}{T}$$



 D_f is increased by

1. decreasing δQ

2. increasing v_0

One way to increase v_0 is to increase $\frac{\Delta p}{p}$. To further increase it we use RF phase displacement, using a high frequency RF cavity. In this case RF buckets are centered on the resonance.



The buckets are empty and beam is forced between them. Now,

$$D_f = \frac{1}{1 + \frac{RB\rho T}{V\frac{\Delta p}{p}} (\omega \delta Q)^2}$$

Without RF phase displacement, a 100 % modulated spill has $D_f = 0.67$. In this case,

$$\omega \delta Q \ge \frac{\Delta p}{p} \frac{1}{T}$$





