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The study of the Network based Event Builder

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Abstract

On 2th June 1999, BELLE detector observed first Beam interactions. The aim of the BELLE experiment is to measure angles of the unitarity triangle and to see the Kobayashi-Masukawa theory which is the origin of CP violation.

The BELLE Event Builder completed on Junuary 1999, and works stablily in efforts of Yamagata-san and members of KEK online group. The role of Event Builder is to collect event records from sub systems and reconstructs full event records with 15MB/s data rate.

In this summer, KEK-B Accelerator acheived a luminosity $10^{32}cm^{-2}s^{-1}$, and we will run at $10^{34}cm^{-2}s^{-1}$ (bunch crossing rate: 500Hz) in the future. In order to endure such high rates, BELLE Event Builder adopts Barrel-Shifter type, but the performance is not enough yet. Now that Evnet Builder works stablily at $10^{32}cm^{-2}s^{-1}$, we should prepre for coming $10^{34}cm^{-2}s^{-1}$ beam.

In this paper, I report the states of Event Builder and my study about Network based type which is one of prominent future plans. To establish the way to construct the event builder using commercial network devices enables us to construct cheap, flexible and reliable DAQ system. Though there are many points to study, I can see the possibility of the network based Event Builder which can have enough performances.

Contents

1 introducion **7**

2 BELLE experiment overview **9**
2.1 Scientific Motivation 9
2.1.1 KM matrix 9
2.1.2 Unitarity triangle 10
2.2 CP violation in B Meson system 11
2.2.1 Direct CP Violation 11
2.2.2 Indirect CP Violation 12
2.2.3 $B^0 \to J/\Psi K_s^0$ 12
2.2.4 $B^0 \to \pi^+\pi^-$ 12
2.3 BELLE experiment 15
2.3.1 KEKB Accerelator 15
2.3.2 SVD(Sillicon Vertex Detector) 17
2.3.3 CDC(Central Drift Chamber) 19
2.3.4 ACC(Aerogel Cerenkov Counters) 20
2.3.5 TOF/TSC 20
2.3.6 ECL(Electromagnetic Calorimetry) 20
2.3.7 KLM(K_L and Muon detector) 22
2.4 Trigger system 22
2.5 Data acquisition system 26
2.5.1 Detector Subsystem 27
2.5.2 Event Builder 27
2.5.3 Online Computer Farm 29
2.5.4 Mass storage system 29
2.5.5 Other subsystem 30

3 BELLE Event Builder **31**
3.1 BELLE event builder over view 31
3.1.1 the switching system 33
3.1.2 Sbus G-Link module TX and RX 34
3.1.3 ECM(External Controll Module) 34

3.2 performance . . . 34
3.2.1 data transfer subsystem to event builder . . . 35
3.2.2 Sbus mmap and Switching . . . 35
3.2.3 summary . . . 38
3.3 problem . . . 39
3.3.1 Scalability . . . 39
3.3.2 GLINK . . . 39
3.3.3 performance . . . 41
3.4 summary . . . 41

4 The network based event builder **42**
4.1 overview . . . 42
4.1.1 setup . . . 42
4.1.2 merits . . . 43
4.2 problems in development . . . 43
4.2.1 software . . . 44
4.2.2 Selection of devices . . . 44
4.2.3 tuning . . . 44
4.2.4 influence of using multi nodes . . . 44
4.3 goals . . . 45

5 Tools and component **46**
5.1 KONOE . . . 46
5.2 DBS . . . 47
5.3 TCP/IP . . . 48
5.4 Fast Ethernet(100BASE-TX) . . . 48
5.5 Switching Hub . . . 48
5.5.1 mechanism of transmission . . . 48
5.5.2 stability of transmitting . . . 49

6 Preparing for the study **52**
6.1 Parameters . . . 52
6.1.1 TCP socket parameter RCV_BUF . . . 52
6.1.2 packet size . . . 54
6.2 Data Reconstruction . . . 55
6.2.1 packet structure . . . 55
6.2.2 packet reconstruction . . . 55
6.3 influence of data reconstruction . . . 55
6.3.1 1×1 test . . . 58
6.3.2 1×4 test . . . 58
6.3.3 the influence of data size on 100 Base-TX . . . 58

7 The study of multipoint network **61**
7.1 The expansion of 1×2 type . . . 61
7.1.1 2×2 setup . . . 61
7.1.2 Considerations . . . 62
7.2 The test of expanded network . . . 63
7.2.1 The expectation of the throughput . . . 63
7.2.2 The expanded network test . . . 64
7.3 Synchronized data transmittion . . . 64
7.3.1 How to synchronize? . . . 65
7.3.2 Measurment of basic performance at 1×1 setup . . . 65
7.3.3 The measurment of multi point setup . . . 67
7.3.4 Expecting throughput . . . 69
7.3.5 Synchronized test using SPARC-5V . . . 69
7.4 Summary . . . 69

8 The possibility of the network based event builder **71**
8.1 Application for BELLE experiment . . . 71
8.1.1 subsystem . . . 71
8.1.2 EB-N . . . 72
8.1.3 timing matrix of event sending . . . 72
8.1.4 summary . . . 74
8.2 enlargement . . . 74
8.2.1 the number of hosts . . . 74
8.2.2 the influence of basic performance . . . 74

9 summary and Future Prospects **75**

A switch test **76**

B data reconstruction **77**

List of Figures

2.1 Unitarity triangle . . . 10
2.2 The tree diagram of $b \to ccs$ . . . 13
2.3 The penguin diagram of $b \to ccs$ . . . 13
2.4 The penguin diagram of $b \to u\bar{u}s$ . . . 14
2.5 The penguin diagram of $b \to u\bar{u}s$ . . . 14
2.6 Tagging and proper time measurement . . . 15
2.7 BELLE detector . . . 16
2.8 Configuration of the KEKB accelerator system . . . 18
2.9 The r-ϕ view of the Silicon Vertex Detector . . . 19
2.10 The r-z view of the Silicon Vertex Detector . . . 19
2.11 Structure of Central drift Chamber . . . 20
2.12 The Aerogel Cherenkov Counter . . . 21
2.13 Time of Flight system . . . 21
2.14 The Electro-Magnetic Calorimeter . . . 23
2.15 Front view of a quadrant the endcap KLM . . . 24
2.16 Front view of a barrel the endcap KLM . . . 24
2.17 Schematic view of trigger system . . . 25
2.18 Global data flow . . . 26
2.19 The architecture of subsystems . . . 27
2.20 The architecture of event builder . . . 28
2.21 The architecture of the online computer farm . . . 29

3.1 usual event builder : only one node is active simultaneously . . . 32
3.2 multiple nodes event builder : reconstruction is done in parallel . . . 32
3.3 4x4 barrel shifter event builder . . . 33
3.4 switching mode . . . 34
3.5 The block diagram of G-Link TX module . . . 35
3.6 The ECM controol timing . . . 36
3.7 The data flow of event builder . . . 37
3.8 The data flow of event builder . . . 37
3.9 The data flow of event builder . . . 38
3.10 The data flow of event builder . . . 38
3.11 The data transfer ability . . . 39

3.12 12×6 switch configuration 40

4.1 The setup of the ethernet based event builder 43

5.1 DBS command file 47
5.2 the movement of Switching Hub : Multipul data streams can go through the switch without interactions 49
5.3 Ethernet flame and transmit starting point 50
5.4 2x1 setup 50
5.5 2x1 throughput : router 51
5.6 2x1 throughput : sw 51

6.1 1x1 setup 52
6.2 socket parameters 53
6.3 RCV_BUF test 53
6.4 packet size test 54
6.5 the structure of the packet 56
6.6 the scheme of packet reconstruction 57
6.7 1×1 test 57
6.8 1×4 test 59
6.9 the influence of the number of event data on throughput 59

7.1 2×2 test 62
7.2 ideal case:synchronized transmittion at 2×2 setup 62
7.3 packets mixing 63
7.4 1×1 + 1×2 test 64
7.5 the synchronized data transmittion 66
7.6 Synchronized 1×1 test 66
7.7 the influence of the number of event data on synchronized 1×1 test 67
7.8 Timing table of 3×2 test: connections of one host are active twice, and non-active once . 68

8.1 Timing matrix of 12×12 setup 73

List of Tables

2.1 Parameters of KEKB . . . 17
2.2 the summary of multi point test . . . 24
2.3 the result of 2×2 test . . . 28

6.1 the result of 1×1 test . . . 58
6.2 the result of 1×4 test . . . 58

7.1 the result of 2×2 test . . . 61
7.2 the result of 1×1 + 1×2 test . . . 63
7.3 the result of the expanded test . . . 65
7.4 the result of synchronized multi point test . . . 68
7.5 the summary of multi point test . . . 70

8.1 the required data size of subdetectors . . . 71

Chapter 1

introducion

Now, the current universe is composed entirely of matter, but it is known that it has been evolved from the matter-antimatter Big-Bang through exceedingly small asymmetry between matter and anti-matter. We could see so-called "CP violation" effects only in nuetral K meson system, but the Standard Model predicted that we can observe larger CP violation in B meson system. The origin of CP violation in the Standard Model is delived from the Kobayashi-Masukawa theory, so we can check the Kobayashi-Masukawa theory by measuring B mesons and determing the value of Kobayashi-Masukawa Matrix.

The Standard Model predicted that CP violation in B meson system is larger thanK meson, but the typical branching ratio of the decay which is usable for the obsevation of CP violation is rather small($10^{-4} \sim 10^{-6}$). In order to generate B mesons in large quantities, an asymmetric e^+e^- collider KEKB was planed, which will run with full bunch (bunch crossing rate: 500MHz) and accomplish the luminosity $\sim 10^{34} cm^{-2} s^{-1}$ and generate 10^8 B mesons a year.

To catch B mesons effectively in the luminosity $10^{34} cm^{-2} s^{-1}$ (1000 times higher than Tristan), we need the excellent triger system and the powerful DAQ system. Since the triger rate and data size are expected as $\sim 500Hz$ (physics event $\sim 100Hz$) and 30KByte/event respectively, the required performances for DAQ system is almost 15MByte/s.

In BELLE experiment, the Event Builder, which constructs full event records, consists of multiple event builder nodes (EB-N) and operates them in parallel, because 15MByte/s data rate exceeds the processing speed available in single node. The processing speed in one EB-N is a few MByte/s, so the 15MByte/s bandwidth is acheived with more than 6 EB-N. 12 sub detector node (IB-N) and EB-N are connected with barrel sifter type switching network and its control module.

When BELLE Event Builder was designed, the switching network which satisfied the requirement of BELLE DAQ system was only barrel sifter types which was developed at Fermilab. Nowadays, the wider computing network spread, the charper and the more powerful Ethernet devices become. It was known that it is possible that we can construct the event builder using Ethernet in principle which is useful, cheap and available. Since there are still many problems to construct the event builder by ethernet, we need to begin studies.

The BELLE experiment began on June 1999, and KEKB acheived the luminosity $10^{32} cm^{-2} s^{-1}$ this summer and aim $10^{33} cm^{-2} s^{-1}$ in autumn 1999. Since the enviroment of DAQ system is expected to be harder and header, we have to check the performance and problems of Event Builder again before the liminorsity is achieved $10^{34} cm^{-2} s^{-1}$.

In this papar, the present status of Event builder and its problems and studies of the ethernet based event builder are reported. Also, the possiblity of the ethernet based event builder is described.

Chapter 2

BELLE experiment overview

The aim of the BELLE experiment is to measure CP violation in B meson system. In this chapter, the process of the B meson causing CP violation and how the BELLE detector will catch it are shown.

2.1 Scientific Motivation

In Standard Model, quark mixing in weak interaction is shown by Kobayashi-Masukawa Model and CP violation is the consquence of a imaginary part of KM matrix.

2.1.1 KM matrix

The relation of quark mass eigenstate and weak eigenstate is shown by unitary matrix, and it has a imaginary phase and called KM matrix when the quark consists of 3 generations.

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = V \begin{pmatrix} d \\ s \\ b \end{pmatrix} \tag{2.1}$$

$$V = \begin{pmatrix} V^{ud} & V^{us} & V^{ub} \\ V^{cd} & V^{cs} & V^{cb} \\ V^{td} & V^{ts} & V^{tb} \end{pmatrix} \tag{2.2}$$

$$= \begin{pmatrix} c_{12} & s_{12}c_{13} & s_{13}e^{-\delta_{13}} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{\delta_{13}} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{\delta_{13}} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{\delta_{13}} & -c_{12}c_{23} - s_{12}c_{23}s_{13}e^{\delta_{13}} & c_{23}c_{13} \end{pmatrix} \tag{2.3}$$

where $s_{ij} = \sin\theta_{ij}$, $c_{ij} = \cos\theta_{ij}$ and δ are free parameters of unitary matrix and i,j $\in$ 1,2,3 are generation labels.

Applying the approximate parameterizaion in (2.3), $s_{12} = \lambda \simeq 0.22$, $s_{23} = A\lambda^2$ and $s_{13}e^{-i\delta} = A\lambda^3(\rho - i\eta)$, this is the Wolfenstein Parameterization which is often used in B physiscs.

$$V \simeq \begin{pmatrix} 1-\lambda^2/2 & \lambda & \lambda^3 A(\rho - i\eta) \\ -\lambda & 1-\lambda^2/2 & \lambda^2 A \\ \lambda^3 A(1-\rho-i\eta) & -\lambda^2 A & 1 \end{pmatrix} + O(\lambda^4) \tag{2.4}$$

2.1.2 Unitarity triangle

The relation of Unitarity of KM matrix shown as

$$V^*_{ij}V^*_{ik} = \delta_{jk} \tag{2.5}$$

can be represented as a triangle in imaginary plane.

In case of nuetral B meson($j = b, k = d$), this triangle is usable because each lengths of sides is same order. Using wolfenstein parameterization,

$$V^*_{tb}V^*_{td}, \quad V^*_{cb}V^*_{cd}, \quad V^*_{ub}V^*_{us} \sim \lambda^3 \tag{2.6}$$

the triangle on imaginary plane, called Unitarity triangle shown in Figure 2.1, has areas ,and three angles are defined as,

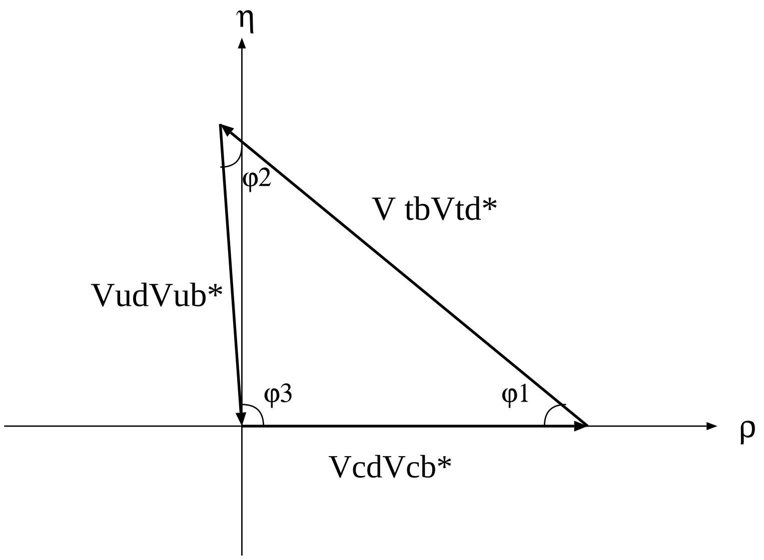

Figure 2.1: Unitarity triangle

$$\phi_1 \equiv arg\Big(\frac{V_{cb}^* V_{cd}}{V_{tb}^* V_{td}}\Big)$$

$$\phi_2 \equiv arg\Big(\frac{V_{ub}^* V_{ud}}{V_{tb}^* V_{td}}\Big)$$

$$\phi_1 \equiv arg\Big(\frac{V_{cb}^* V_{cd}}{V_{ub}^* V_{ud}}\Big)$$

In oreder to test the KM model as a origin of CP violation, it's very effective to measure the length of sides and angles of the Unitarity triangle. If angels are small and the triangle has no area, there would be a new physics beyond the Standard model.

2.2 CP violation in B Meson system

2.2.1 Direct CP Violation

CP symmetry is violated when there is a complex coupling constant. We can take a whole phase factor arbitrarily, but when we cannot make all of the coupling constants real, CP violation occurs.

Direct CP violation can be observerd as the difference of decay probability between one mode and its CP transfered mode only when there are more than two independent amplitude contribution to the probability. Consider the mode $X \to Y$ which has two independent amplitude.

$$\begin{aligned} Amp(X \to Y) &= A_1 e^{i\delta_1} e^{i\phi_1} + A_2 e^{i\delta_2} e^{i\phi_2} \\ &= A_1 e^{i\delta_1} e^{i\phi_1} \left[1 + \frac{A_2}{A_1} e^{i(\delta_2-\delta_1)} e^{i(\phi_2-\phi_1)}\right], \end{aligned} \tag{2.7}$$

where $\delta_{1,2}$ are the strong phase which are independent of weak phase $\phi_{1,2}$. The decay probability is proportional to square of the absolute value of the amplitude.

$$|Amp(X \to Y)|^2 = A_1^2 \left[1 + \frac{A_2^2}{A_1^2} + 2\frac{A_2}{A_1}\cos(\delta_2 - \delta_1 + \phi_2 - \phi_1)\right]$$

For the CP transformed mode, weak phase changes its sign and strong phase does not.

$$|Amp(X^{CP} \to Y^{CP})|^2 = A_1^2 \left[1 + \frac{A_2^2}{A_1^2} + 2\frac{A_2}{A_1}\cos(\delta_2 - \delta_1 - \phi_2 + \phi_1)\right].$$

So, the asymmetry arises,

$$Asym \propto |Amp(X \to Y)|^2 - |Amp(X^{CP} \to Y^{CP})|^2 = -4A_2A_1 \sin(\delta_2 - \delta_1)\sin(\phi_2 - \phi_1).$$

Thus, to observe the violation, the mode has to have two independent amplitude, and both the weak phase and strong phase must be different between the two amplitude ($\delta_2 \neq \delta_1, \;\; \phi_2 \neq \phi_1$).

2.2.2 Indirect CP Violation

When B^0 and $\bar{B}^0$ mesons decay into the same CP eigenstate f, the CP violation aasymmtry A_f is given by,

$$A_f = \frac{R(B^0 \to f) - R(\bar{B^0} \to f)}{R(B^0 \to f) + R(\bar{B^0} \to f)} = sin(2\phi_{CP})\ sin(\Delta m \Delta t) \tag{2.8}$$

where Δm is the mass defference between two B^0's and Δt is the difference of proper time for B^0 and $\bar{B^0}$ decays. $R(B^0 \to f)$ is the decay rate of $B^0 \to f$.

The asymmetry is produced by the interference of two amplitudes; a direct decay $B^0 \to f$ and a decay through the $B\bar{B}$ 0scillation $B^0 \to \bar{B}^0 \to f$ which is called as the indirect CP violation. Experimentally, the measurement of Δt is essential for the observation of the asymmetry. We can measure Δt using the Lorentz boost of the $\Upsilon(4S)$ state produced by the asymmetry of beam energies as described later. The angle ϕ_{CP} is the phase difference between two amplitude, which is derectly related to the internal angle of the unitarity triangle.

In BELLE experiment, we try to measure angles of the unitarity triangle by following decays to CP eigenstate.

$$\phi_1 : B^0 \to J/\Psi K_s^0$$
$$\phi_2 : B^0 \to \pi^+\pi^-$$
$$\phi_3 : B^0 \to D^0 K$$

2.2.3 $B^0 \to J/\Psi K_s^0$

The decay mode $b \to ccs$ is used to measure the angle ϕ_1. The time dependent CP asymmetry is written as follows;

$$Asym[J/\Psi K_s^0] = sin(2\phi_1)\ sin(\Delta m \Delta t)$$

There're two large contrubutions for this decay; one is tree diagram (Figure 2.2) and the other is pengin diagram (Figure 2.3).

The contribution in weak phase from the penguin diagram is written as the sum of three terms, $(V_{ub}^* V_{us} + V_{cb}^* V_{cs} + V_{tb}^* V_{ts})$, which can be expressed as $\sim V_{cb}^* V_{cs}$, using the unitarity relation $V_{ib}^* V_{is} = 0$. So the weak phase penguin diagram is approximately same as $V_{cb}^* V_{cs}$, which is same with that of tree diagram.

There're two large contrubutions for this decay, but weak phases are almost same and direct CP violation hardly occures. Also, $J/\Psi \to l^+ l^-$ and $K_s^0 \to \pi^+\pi^-$ is very clean. This is why the decay mode $B^0 \to J/\Psi K_s^0$ is called as "Gold Plate Mode".

2.2.4 $B^0 \to \pi^+\pi^-$

The decay mode $b \to u\bar{u}d$ is used to measure the angle ϕ_2. The time dependent CP asymmetry is, if the contribution is only from tree diagram (Figure 2.4),

$$Asym[\pi^+\pi^-] = sin(2\phi_2)\ sin(\Delta m \Delta t)$$